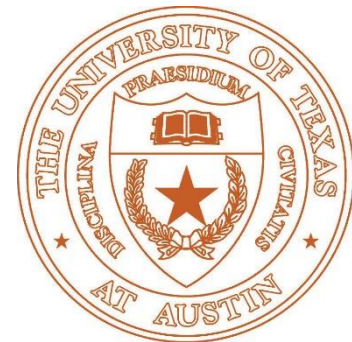
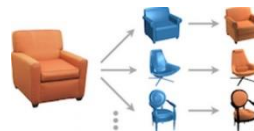
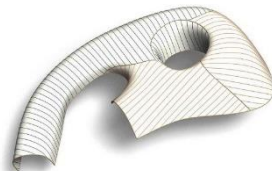
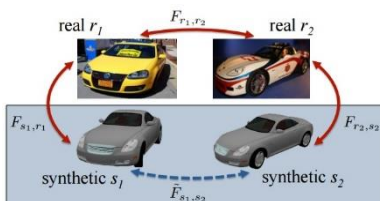
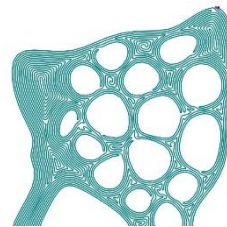


CS354 Computer Graphics

Particle Systems

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Reading

- Required:
 - Witkin, *Particle System Dynamics*, SIGGRAPH '97 course notes on Physically Based Modeling.
 - Witkin and Baraff, *Differential Equation Basics*, SIGGRAPH '01 course notes on Physically Based Modeling.
- Optional
 - Hockney and Eastwood. *Computer simulation using particles*. Adam Hilger, New York, 1988.
 - Gavin Miller. "The motion dynamics of snakes and worms." *Computer Graphics* 22:169-178, 1988.

What are particle systems?

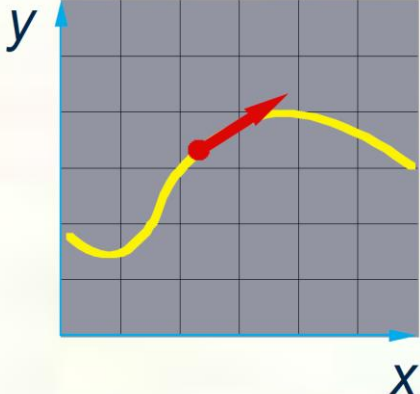
- A **particle system** is a collection of point masses that obeys some physical laws (e.g, gravity, heat convection, spring behaviors, ...).
- Particle systems can be used to simulate all sorts of physical phenomena:

Particle in a flow field

- We begin with a single particle with:

■ Position, $\vec{\mathbf{x}} = \begin{bmatrix} x \\ y \end{bmatrix}$

■ Velocity, $\vec{\mathbf{v}} = \dot{\vec{\mathbf{x}}} = \frac{d\vec{\mathbf{x}}}{dt} = \begin{bmatrix} dx/dt \\ dy/dt \end{bmatrix}$

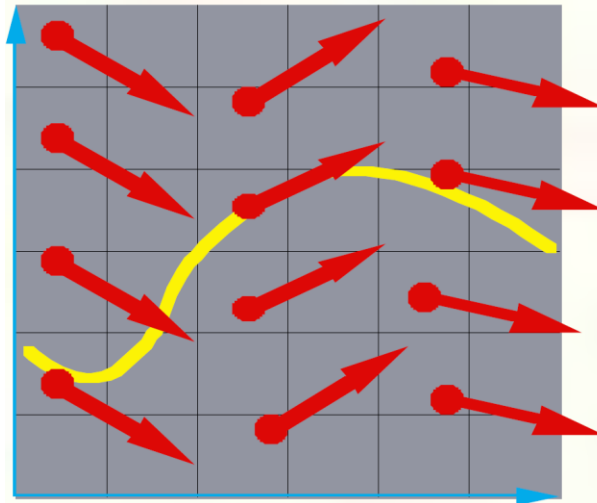


The diagram shows a 2D coordinate system with a grid. The horizontal axis is labeled 'x' and the vertical axis is labeled 'y'. A yellow curve represents the trajectory of a particle, starting from the left, moving up and then curving to the right. A red arrow is drawn tangent to the curve at a specific point, representing the velocity vector at that position.

- Suppose the velocity is actually dictated by some driving function \mathbf{g} : $\dot{\vec{\mathbf{x}}} = \mathbf{g}(\vec{\mathbf{x}}, t)$

Vector fields

- At any moment in time, the function \mathbf{g} defines a vector field over \mathbf{x} :



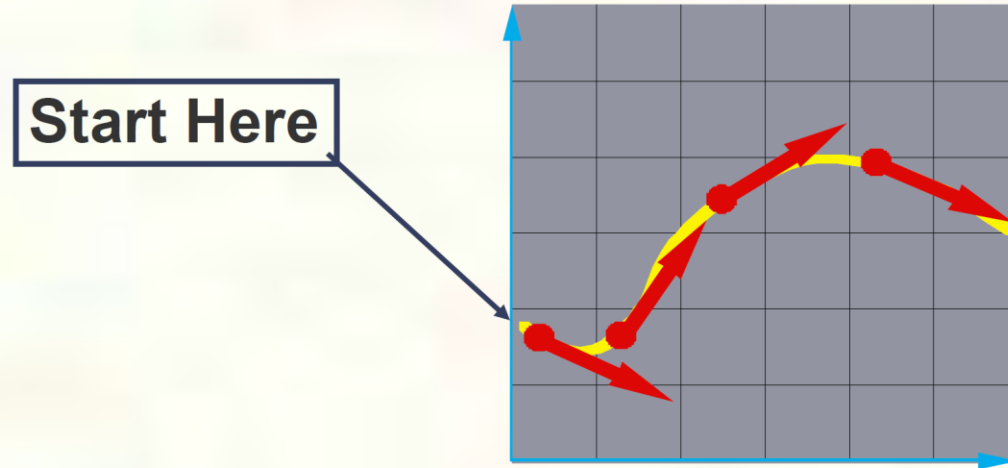
- How does our particle move through the vector field?

Diff eqs and integral curves

- The equation $\dot{\mathbf{x}} = g(\vec{\mathbf{x}}, t)$

is actually a **first order differential equation**.

- We can solve for \mathbf{x} through time by starting at an initial point and stepping along the vector field:



- This is called an **initial value problem** and the solution is called an **integral curve**.

Euler's method

- One simple approach is to choose a time step, Δt , and take linear steps along the flow:
$$\vec{\mathbf{x}}(t + \Delta t) = \vec{\mathbf{x}}(t) + \Delta t \cdot \dot{\vec{\mathbf{x}}}(t) = \vec{\mathbf{x}}(t) + \Delta t \cdot g(\vec{\mathbf{x}}, t)$$

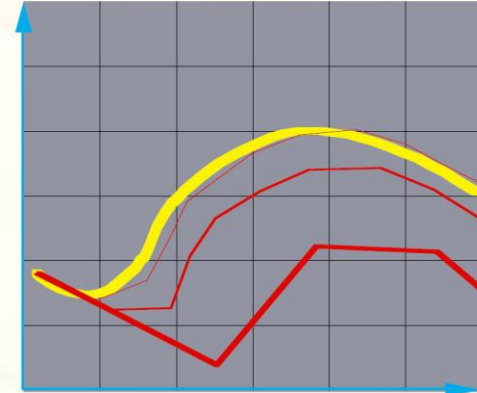
- Writing as a time iteration:
$$\vec{\mathbf{x}}^{i+1} = \vec{\mathbf{x}}^i + \Delta t \cdot \vec{\mathbf{v}}^i$$

- This approach is called **Euler's method** and looks like:

- Properties:

- Simplest numerical method
- Bigger steps, bigger errors. Error $\sim O(\Delta t^2)$.

- Need to take pretty small steps, so not very efficient. Better (more complicated) methods exist, e.g., “Runge-Kutta” and “implicit integration.”



Particle in a force field

- Now consider a particle in a force field \mathbf{f} .
- In this case, the particle has:
 - Mass, m
 - Acceleration, $\vec{\mathbf{a}} \equiv \ddot{\mathbf{x}} = \frac{d\vec{\mathbf{v}}}{dt} = \frac{d^2\vec{\mathbf{x}}}{dt^2}$
- The particle obeys Newton's law: $\vec{\mathbf{f}} = m\vec{\mathbf{a}} = m\ddot{\mathbf{x}}$
- The force field \mathbf{f} can in general depend on the position and velocity of the particle as well as time.
- Thus, with some rearrangement, we end up with:

$$\ddot{\mathbf{x}} = \frac{\vec{\mathbf{f}}(\vec{\mathbf{x}}, \dot{\mathbf{x}}, t)}{m}$$

Second order equations

This equation:

$$\ddot{\vec{x}} = \frac{\vec{f}(\vec{x}, \dot{\vec{x}}, t)}{m}$$

is a **second order differential equation**.

Our solution method, though, worked on first order differential equations.

We can rewrite this as:

$$\begin{bmatrix} \dot{\vec{x}} = \vec{v} \\ \dot{\vec{v}} = \frac{\vec{f}(\vec{x}, \vec{v}, t)}{m} \end{bmatrix}$$

where we have added a new variable **v** to get a pair of coupled first order equations.

Phase space

$$\begin{bmatrix} \vec{x} \\ \vec{v} \end{bmatrix}$$

- Concatenate \mathbf{x} and \mathbf{v} to make a 6-vector: position in **phase space**.

$$\begin{bmatrix} \dot{\vec{x}} \\ \dot{\vec{v}} \end{bmatrix}$$

- Taking the time derivative: another 6-vector.

$$\begin{bmatrix} \dot{\vec{x}} \\ \dot{\vec{v}} \end{bmatrix} = \begin{bmatrix} \vec{v} \\ \vec{f}/m \end{bmatrix}$$

- A vanilla 1st-order differential equation.

Differential equation solver

Starting with:

$$\begin{bmatrix} \dot{\vec{x}} \\ \dot{\vec{v}} \end{bmatrix} = \begin{bmatrix} \vec{v} \\ \vec{f}/m \end{bmatrix}$$

Applying Euler's method:

$$\vec{x}(t + \Delta t) = \vec{x}(t) + \Delta t \cdot \dot{\vec{x}}(t)$$

$$\dot{\vec{x}}(t + \Delta t) = \dot{\vec{x}}(t) + \Delta t \cdot \ddot{\vec{x}}(t)$$

And making substitutions:

$$\vec{x}(t + \Delta t) = \vec{x}(t) + \Delta t \cdot \vec{v}(t)$$

$$\dot{\vec{x}}(t + \Delta t) = \dot{\vec{x}}(t) + \Delta t \cdot \vec{f}(\vec{x}, \dot{\vec{x}}, t) / m$$

Writing this as an iteration, we have:

$$\vec{x}^{i+1} = \vec{x}^i + \Delta t \cdot \vec{v}^i$$

$$\vec{v}^{i+1} = \vec{v}^i + \Delta t \cdot \frac{\vec{f}^i}{m}$$

Again, performs poorly for large Δt .

Verlet Integration

- Also called Størmer's Method
 - Invented by Delambre (1791), Størmer (1907), Cowell and Crommelin (1909), Verlet (1960) and probably others
- More stable than Euler's method (time reversible as well)

Forces

- Each particle can experience a force which sends it on its merry way.
- Where do these forces come from? Some examples:
 - Constant (gravity)
 - Position/time dependent (force fields)
 - Velocity-dependent (drag)
 - Combinations (Damped springs)
- How do we compute the net force on a particle?

Gravity and viscous drag

The force due to **gravity** is simply:

$$\vec{\mathbf{f}}_{grav} = m\vec{\mathbf{G}}$$

$$\text{p->f} += \text{p->m} * \text{F->G}$$

Often, we want to slow things down with **viscous drag**:

$$\vec{\mathbf{f}}_{drag} = -k\vec{\mathbf{v}}$$

$$\text{p->f} -= \text{F->k} * \text{p->v}$$

Damped spring

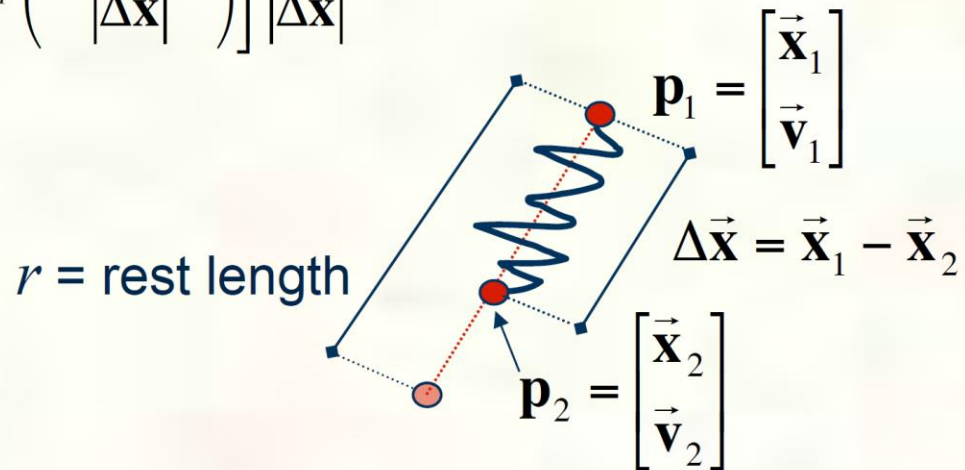
Recall the equation for the force due to a spring: $f = -k_{spring} (|\Delta\vec{\mathbf{x}}| - r)$

We can augment this with damping: $f = -\left[k_{spring} (|\Delta\vec{\mathbf{x}}| - r) + k_{damp} |\vec{\mathbf{v}}|\right]$

The resulting force equations for a spring between two particles become:

$$\vec{\mathbf{f}}_1 = -\left[k_{spring} (|\Delta\vec{\mathbf{x}}| - r) + k_{damp} \left(\frac{\Delta\vec{\mathbf{v}} \cdot \Delta\vec{\mathbf{x}}}{|\Delta\vec{\mathbf{x}}|}\right)\right] \frac{\Delta\vec{\mathbf{x}}}{|\Delta\vec{\mathbf{x}}|}$$

$$\vec{\mathbf{f}}_2 = -\vec{\mathbf{f}}_1$$



derivEval

Clear forces

Loop over particles,
zero force
accumulators

Calculate forces

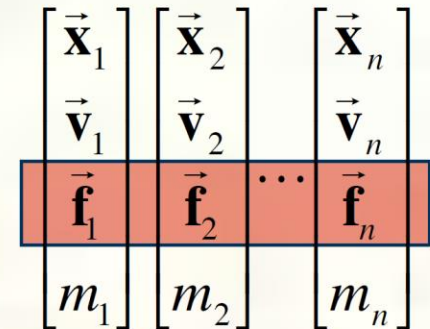
Sum all forces into
accumulators

Return derivatives

Loop over particles,
return \mathbf{v} and \mathbf{f}/m

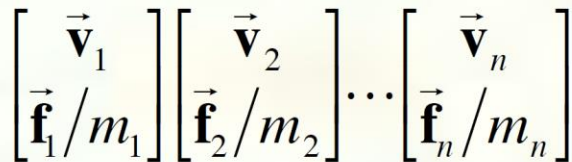
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Clear force
accumulators



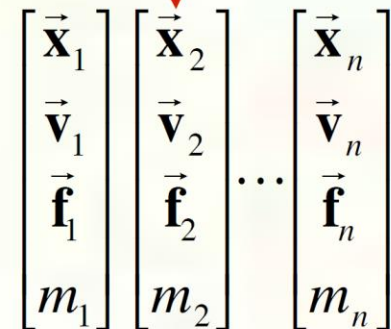
Apply forces
to particles

2

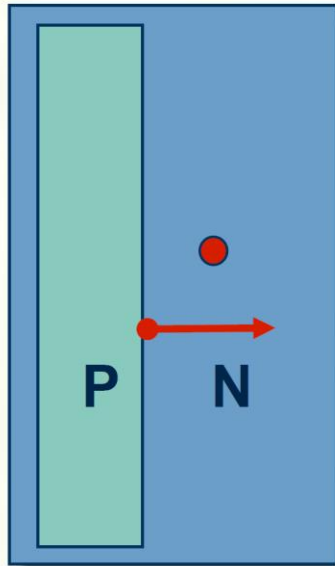


Return derivatives
to solver

3



Bouncing off the walls

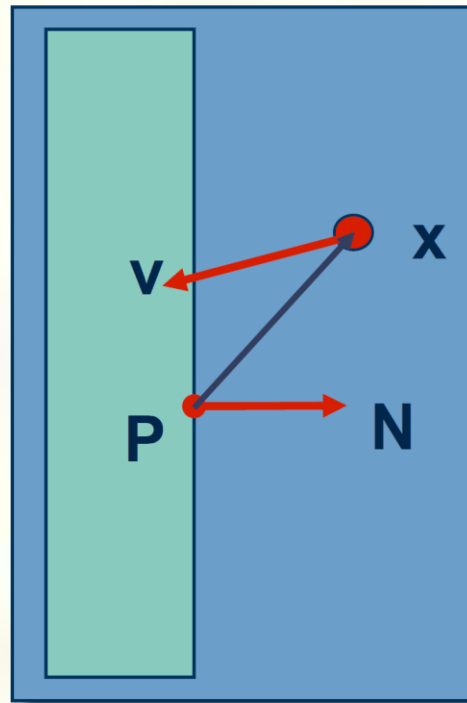


- Add-on for a particle simulator
- For now, just simple point-plane collisions

A plane is fully specified by any point P on the plane and its normal N .

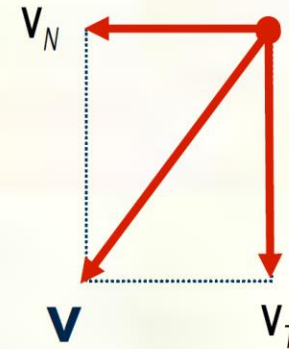
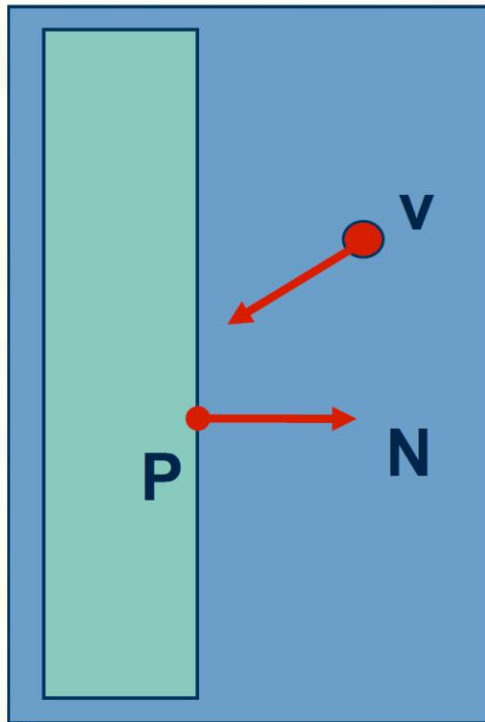
Collision Detection

How do you decide when you've crossed a plane?



Normal and tangential velocity

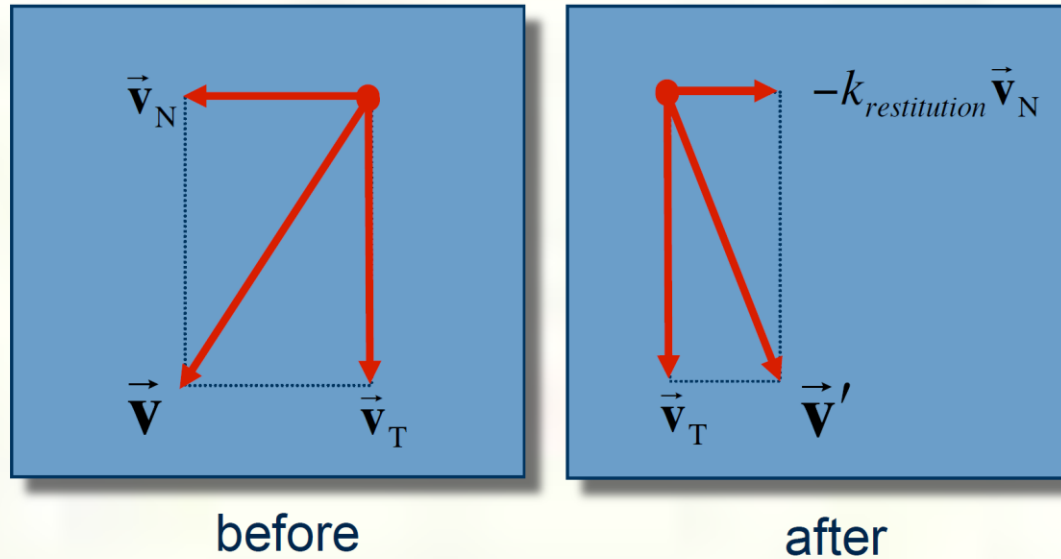
To compute the collision response, we need to consider the normal and tangential components of a particle's velocity.



$$\vec{\mathbf{v}}_N = (\vec{\mathbf{N}} \cdot \vec{\mathbf{v}}) \vec{\mathbf{N}}$$

$$\vec{\mathbf{v}}_T = \vec{\mathbf{v}} - \vec{\mathbf{v}}_N$$

Collision Response



$$\vec{V}' = \vec{V}_T - k_{restitution} \vec{V}_N$$

Without backtracking, the response may not be enough to bring a particle to the other side of a wall.

In that case, detection should include a velocity check:

Discussion