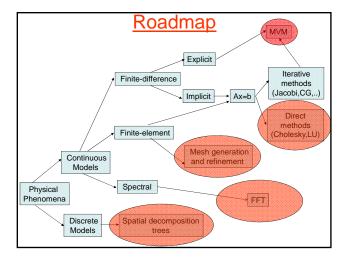
Some Computational Science **Algorithms**

Computational science

- Simulations of physical phenomena
 - fluid flow over aircraft (Boeing 777)fatigue fracture in aircraft bodies

 - evolution of galaxies
- Two main approaches
 - continuous models: fields and differential equations (eg. Navier-Stokes equations, Maxwell's equations,...)
 - discrete models: particles and forces (eg. gravitational forces)
- Paradox
 - most differential equations cannot be solved exactly
 - must use numerical techniques that convert calculus problem to matrix computations: discretization

 - n-body methods are straight-forward
 but need to use a lot of bodies to get accuracy
 - must find a way to reduce O(N2) complexity of obvious algorithm



Organization

- Finite-difference methods
 - ordinary and partial differential equations
 discretization techniques
 explicit methods: Forward-Euler method
 implicit methods: Backward-Euler method
- Finite-element methods

 mesh generation and refinement
- weighted residuals
 N-body methods Barnes-Hut
- Key algorithms and data structures

 - Key algorithms and data structures

 matrix computations

 algorithms

 solution of systems of linear equations

 direct methods

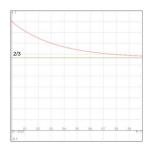
 data structures

 dense and sparse matrices

 graph computations
 - mesh generation and refinement
 spatial decomposition trees

Ordinary differential equations

- · Consider the ode u'(t) = -3u(t)+2u(0) = 1
- This is called an initial value problem
 - initial value of u is given - compute how function u evolves for t > 0
- Using elementary calculus, we can solve this ode exactly u(t) = 1/3 (e^{-3t}+2)

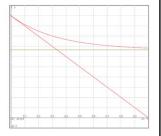


Problem

- For general ode's, we may not be able to express solution in terms of elementary functions
- In most practical situations, we do not need exact solution anyway
 - enough to compute an approximate solution, provided
 - · we have some idea of how much error was introduced
 - we can improve the accuracy as needed
- · General solution:
 - convert calculus problem into algebra/arithmetic problem
 - discretization: replace continuous variables with discrete variables
 - · in finite differences.
 - time will advance in fixed-size steps: t=0,h,2h,3h,...
 - differential equation is replaced by difference equation

Forward-Euler method

- Intuition:
 - we can compute the derivative at t=0 from the differential equation u'(t) = -3u(t)+2
 - so compute the derivative at t=0 and advance along tangent to t =h to find an approximation to u(h)
- Formally, we replace derivative with forward difference to get a difference equation
- u'(t) → (u(t+h) u(t))/h
 Replacing derivative with difference is essentially the inverse of how derivatives were probably introduced to you in elementary calculus



Back to ode

- · Original ode
 - u'(t) = -3u(t)+2
- · After discretization using Forward-Euler:

 $(u_f(t+h) - u_f(t))/h = -3u_f(t)+2$

- After rearrangement, we get difference equation $u_f(t+h) = (1-3h)u_f(t)+2h$
- · We can now compute values of u:

 $u_f(0) = 1$

 $u_f(h) = (1-h)$

 $u_f(2h) = (1-2h+3h^2)$

Exact solution of difference equation

- · In this particular case, we can actually solve difference equation exactly
- · It is not hard to show that if difference equation is $u_f(t+h) = a^*u_f(t)+b$ $u_{f}(0) = 1$

the solution is

 $u_f(nh) = a^n + b^*(1-a^n)/(1-a)$

• For our difference equation,

 $u_f(t+h) = (1-3h)u_f(t)+2h$

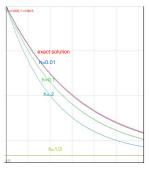
the exact solution is

 $u_f(nh) = 1/3((1-3h)^n+2)$

Comparison

- u(t) = 1/3 (e^{-3t}+2) u(nh) = 1/3(e^{-3nh}+2) (at time-steps) Forward-Euler solution $u_f(nh) = 1/3((1-3h)^n+2)$
- Use series expansion to compare u(nh) = 1/3(1-3nh+9/2 n²h²···+ 2) $u_1(nh) = 1/3(1-3nh+n(n-1)/2 9h^2+...+2)$ So error = O(nh²) (provided h < 2/3)
- Conclusion:
- error per time step (local error) = O(h²)
- U(n*)

 error at time nh = O(nh²)
 In general, Forward-Euler
 converges only if time step is
 "small enough"

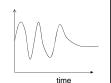


Choosing time step

- Time-step needs to be small enough to capture highest frequency phenomenon of interest
- Nyquist's criterion
 - sampling frequency must be at least twice highest frequency to prevent aliasing
 - for most finite-difference formulas, you need sampling frequencies (much) higher than the Nyquist criterion
- In practice, most functions of interest are not band-limited, so use
 insight from application or
- reduce time-step repeatedly till changes are not significant

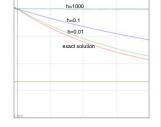
 Fixed-size time-step can be inefficient if frequency varies widely over time interval.

 | The content of the content of
 - other methods like finite-elements permit variable time-steps as we will see later



Backward-Euler method

- Replace derivative with backward difference $u'(t) \rightarrow (u(t) u(t-h))/h$
- For our ode, we get $u_b(t)-u_b(t-h)/h = -3u_b(t)+2$ which after rearrangement $u_b(t)=(2h+u_b(t-h))/(1+3h)$
- As before, this equation is simple enough that we can write down the exact solution:
- $u_b(nh) = ((1/(1+3h))^n + 2)/3$ Using series expansion, we get $u_b(nh) = (1-3nh + (-n(-n-1)/2) 9h^2 + ...+2)/3$

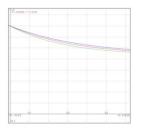


t-h t t+h

Comparison

- Exact solution
 - $u(t) = 1/3 (e^{-3t}+2)$ $u(nh) = 1/3(e^{-3nh}+2)$ (at time-steps) Forward-Euler solution
- $u_f(nh) = 1/3((1-3h)^n+2)$ error = O(nh²) (provided h < 2/3)
- Backward-Euler solution $u_b(n^*h) = 1/3 ((1/(1+3h))^n + 2)$ error = O(nh²) (h can be any value you want)
- Many other discretization schemes have been studied in the literature

 - Runge-KuttaCrank-Nicolson
 - Upwind differencing



Red: exact solution

Blue: Backward-Euler solution (h=0.1) Green: Forward-Euler solution (h=0.1)

Higher-order difference formulas

- First derivatives:
 - Forward-Euler: y'(t) → $y_f(t+h)-y_f(t)$ /h
 - Backward-Euler: y'(t) → $y_b(t)-y_b(t-h)$ /h
 - Centered: $y'(t) \rightarrow y_c(t+h)-y_c(t-h)/2h$
- Second derivatives:
- Forward: y"(t) →

 $(y_f(t+2h)-y_f(t+h))-(y_f(t+h)-y_f(t))/h^2$ $= y_f(t+2h)-2y_f(t+h)+y_f(t)/h^2$

- Backward: $y''(t) \rightarrow y_b(t)-2y_b(t-h)+y_b(t-2h)/h^2$
- Centered: $y''(t) \rightarrow y_c(t+h) 2y_c(t) + y_c(t-h)/h^2$

Systems of ode's

• Consider a system of coupled ode's of the form

$$u'(t) = a_{11}^*u(t) + a_{12}^*v(t) + a_{13}^*w(t) + c_1(t)$$

$$v'(t) = a_{21}^*u(t) + a_{22}^*v(t) + a_{23}^*w(t) + c_2(t)$$

$$w'(t) = a_{31}^*u(t) + a_{32}^*v(t) + a_{33}^*w(t) + c_3(t)$$

• If we use Forward-Euler method to discretize this system, we get the following system of simultaneous equations

$$\begin{split} &u_f(t+h)-u_f(t) \ /h = a_{11}{}^*u_f(t) + a_{12}{}^*v_f(t) + a_{13}{}^*w_f(t) + c_1(t) \\ &v_f(t+h)-v_f(t) \ /h = a_{21}{}^*u_f(t) + a_{22}{}^*v_f(t) + a_{23}{}^*w_f(t) + c_2(t) \\ &w_f(t+h)-w_f(t) \ /h = a_{31}{}^*u_f(t) + a_{32}{}^*v_f(t) + a_{33}{}^*w_f(t) + c_3(t) \end{split}$$

Forward-Euler (contd.)

Rearranging, we get

 $u_f(t+h) = (1+ha_{11})^*u_f(t) + ha_{12}^*v_f(t) + ha_{13}^*w_f(t) + hc_1(t)$ $v_f(t+h) = ha_{21}^*u_f(t) + (1+ha_{22})^*v_f(t) + ha_{23}^*w_f(t) + hc_2(t)$ $w_f(t+h) = ha_{31}^*u_f(t) + ha_{32}^*v_f(t) + (1+a_{33})^*w_f(t) + hc_3(t)$

Introduce vector/matrix notation

 $\underline{U}(t) = [u(t) \ v(t) \ w(t)]^T$

 $C(t) = [c_1(t) c_2(t) c_3(t)]^T$

Vector notation

· Our systems of equations was

 $u_f(t+h) = (1+ha_{11})^*u_f(t) + ha_{12}^*v_f(t) + ha_{13}^*w_f(t) + hc_1(t)$ $\begin{aligned} v_i(t+h) &= ha_{21}^* u_i(t) + (1+ha_{22})^* v_i(t) + ha_{23}^* w_i(t) + hc_2(t) \\ w_i(t+h) &= ha_{31}^* u_i(t) + ha_{32}^* v_i(t) + (1+a_{33})^* w_i(t) + hc_3(t) \end{aligned}$

- This system can be written compactly as follows $\underline{U}(t+h) = (I+hA)\underline{U}(t)+h\underline{C}(t)$
- We can use this form to compute values of $\underline{U}(h),\underline{U}(2h),\underline{U}(3h),...$
- Forward-Euler is an example of explicit method of discretization
 - key operation: matrix-vector (MVM) multiplication
 - in principle, there is a lot of parallelism
 O(n²) multiplications

 - O(n) reductions
 - parallelism is independent of runtime values

Backward-Euler

We can also use Backward-Euler method to discretize system of ode's

 $u_b(t)-u_b(t-h)/h = a_{11}^*u_b(t) + a_{12}^*v_b(t) + a_{13}^*w_b(t) + c_1(t)$ $v_b(t) - v_b(t-h) /h = a_{21}^* v_b(t) + a_{22}^* v_b(t) + a_{23}^* w_b(t) + c_2(t)$ $w_b(t) - w_b(t-h) /h = a_{31}^* u_b(t) + a_{32}^* v_b(t) + a_{33}^* w_b(t) + c_3(t)$

- We can write this in matrix notation as follows $(I\text{-}hA)\underline{U}(t)=\underline{U}(t\text{-}h)\text{+}h\underline{C}(t)$
- Backward-Euler is example of implicit method of discretization
 - key operation: solving a dense linear system $M\underline{x} = \underline{v}$
- · How do we solve large systems of linear equations?
- Matrix (I-hA) is often very sparse
 - Important to exploit sparsity in solving linear systems

Diversion: Solving linear systems

Solving linear systems

- Linear system: $A\underline{x} = \underline{b}$
- Two approaches
 - direct methods: Cholesky, LU with pivoting
 - factorize A into product of lower and upper triangular matrices A =
 - · solve two triangular systems

Ly = b

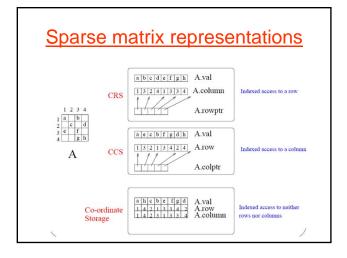
- problems:
 - even if A is sparse, L and U can be quite dense ("fill")
- no useful information is produced until the end of the procedure
- iterative methods: Jacobi, Gauss-Seidel, CG, GMRES
 - guess an initial approximation \underline{x}_0 to solution
 - error is Ax₀ b (called residual)
 - repeatedly compute better approximation \underline{x}_{i+1} from residual $(A\underline{x}_i \underline{b})$
 - · terminate when approximation is "good enough"

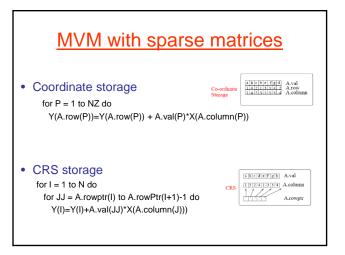
Jacobi iteration: general picture

- Linear system Ax = b
- Jacobi iteration

 $M^{\star}x_{i+1}$ = (M-A)x_i + b (where M is the diagonal of A) This can be written as

- $x_{i+1} = x_i M^{-1}(Ax_i b)$ Key operation:
 - matrix-vector multiplication
- Caveat:
 - Jacobi iteration does not always converge
 - even when it converges, it usually converges slowly
 - there are faster iterative methods available: CG,GMRES,...
 - what is important from our perspective is that key operation in all these iterative methods is matrix-vector multiplication

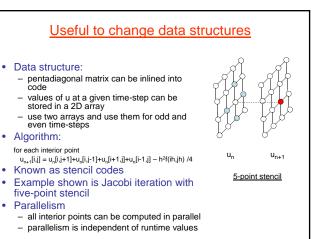




Finite-differences:pde's

Finite-difference methods for solving partial differential equations Basic ideas carry over unchanged Example: 2-d heat equation -'u'-x² + -'u'-y² = f(x,y) assume temperature at boundary is fixed Discretize domain using a regular NxN grid of pitch h Approximate derivatives as centered differences -'u'-y² -> ((u(i,j+1)-u(i,j))/h - (u(i,j)-u(i,j))/h)/h -'u'-x² -> ((u(i+1,j)-u(i,j))/h - (u(i,j)-u(i-1,j))/h)/h So we get a system of (N-1)x(N-1) difference equations in terms of the unknowns at the (N-1)x(N-1) interior points j interior point (i,j) u(i,j+1)+u(i,j-1)+u(i+1,j)+u(i-1,j) - 4u(i,j) = h² f(ih,jh) This system can be solved using any of our methods.

Solving partial differential equations contd.) System of (N-1)x(N-1) difference equations in terms of the unknowns at the (N-1)x(N-1) interior points $u(i,j+1)+u(i,j-1)+u(i+1,j)+u(i-1,j)-4u(i,j)=h^2 f(ih,jh)$. Matrix notation: use row-major (natural) order for u's u(i-1,j) \circ u(i,j-1) 0..10..01-4 10..010...0. $= h^2 f(ih,jh)$ u(i,j) u(i,j+1) 0..0010..01-410..010. 5-point stencil u(i+1,j) Pentadiagonal sparse matrix Can be represented using specialized sparse matrix formats Since matrix is sparse, we should use an iterative method like Jacobi.



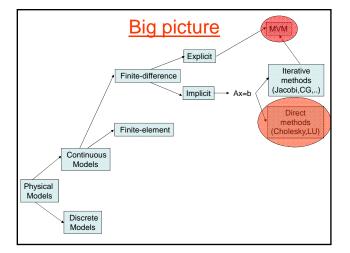
Observations

- Algorithm: Jacobi iteration with 5-point stencil to solve 2-D heat equation
- Two very different programs
 - a. pentadiagonal matrix: stored in sparse matrix format, unknowns: 1D vector
 - b. pentadiagonal matrix: inlined into code, unknowns: matrix
- · Data structures are critical
 - can result in very different programs (implementations) for the same algorithm

Summary

- Finite-difference methods
 - can be used to find approximate solutions to ode's and pde's
- Many large-scale computational science simulations use these methods
- Time step or grid step needs to be constant and is determined by highest-frequency phenomenon
 - can be inefficient for when frequency varies widely in domain of interest
 - one solution: structured AMR methods





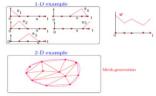
Finite-element methods

- Express approximate solution to pde as a linear combination of certain basis functions
- Similar in spirit to Fourier analysis
 - express periodic functions as linear combinations of sines and cosines
- Questions:
 - what should be the basis functions?

 - mesh generation: discretization step for finite-elements
 mesh defines basis functions "3," 4," 5,....which are low-degree piecewise polymerial functions." polynomial functions
 - given the basis functions, how do we find the best linear combination of these for approximating solution to pde?

 - $\mathbf{u} = \Sigma_1 \, \mathbf{c}_1^{-\epsilon}$ weighted residual method: similar in spirit to what we do in Fourier analysis, but more complex because basis functions are not necessarily orthogonal

Mesh generation and refinement



- - 1-D example:

 mesh is a set of points, not necessarily equally spaced

 basis functions are "hats" which

 have a value of 1 at a mesh points

 deay down to 0 at neighboring mesh points

 0 everywhere else

 linear combinations of these produce piecewise linear functions in domain, which may change slope only at mesh points

 In 2-D, mesh is a triangularization of domain, while in 3-D, it might be a tetrahedralization
- Mesh refinement: called h-refinement

 - add more points to mesh in regions where discretization error is large irregular nature of mesh makes this easy to do this locally finite-differences require global refinement which can be computationally expensive

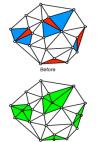
Delaunay Mesh Refinement

Iterative refinement to remove bad triangles with lots of discretization error:

while there are bad triangles do (
Pick a bad triangle;
Find its cavity;
Retriangulate cavity;
// may create new bad triangles

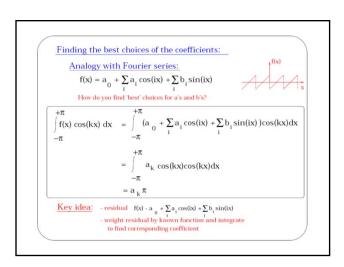
- Don't-care non-determinism:
- final mesh depends on order in which bad triangles are processed applications do not care which mesh is produced
- Data structure:
 - graph in which nodes represent triangles and edges represent triangle adjacencies
 - Parallelism:

 - bad triangles with cavities that do not overlap can be processed in parallel parallelism is dependent on runtime values compilers cannot find this parallelism
 - (Miller et al) at runtime, repeatedly build interference graph and find maximal independent sets for parallel execution

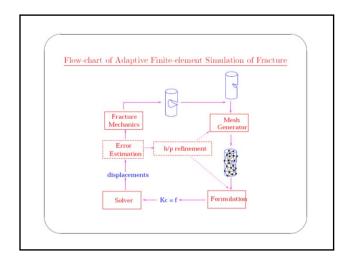


Finding coefficients

- · Weighted residual technique
 - similar in spirit to what we do in Fourier analysis, but basis functions are not necessarily orthogonal
- · Key idea:
 - problem is reduced to solving a system of equations $A\underline{x} = \underline{b}$
 - solution gives the coefficients in the weighted sum
 - because basis functions are zero almost everywhere in the domain, matrix A is usually very sparse
 - number of rows/columns of A ~ O(number of points in mesh)
 - number of non-zeros per row ~ O(connectivity of mesh point)
 - typical numbers:
 - A is 10⁶x10⁶
 - only about ~100 non-zeros per row



```
\begin{array}{ll} & \frac{Weighted\ Residual\ Technique:}{Residual:} & (L\ u^*-f) = (L\ (\sum\limits_{i=1}^{N} c_i \varphi_i) - f) \\ & Weighted\ Residual = (L\ (\sum\limits_{i=1}^{N} c_i \varphi_i) - f) & \varphi_k \\ & Equation\ for\ k^t \ unknown: & \int\limits_{\Omega} \varphi_k^* (L\ (\sum\limits_{i=1}^{N} c_i \varphi_i) - f) \ dV = 0 \\ & \Rightarrow \\ & If\ the\ differential\ equation\ is\ linear: & \\ & c_i \int\limits_{\Omega} \varphi_k^* L \varphi_i \ dV + .... + c_N \int\limits_{\Omega} \varphi_k^* L \varphi_N \ dV & = \int\limits_{\Omega} \varphi_k \ f \ dV \\ & K = 1.2...N \\ & K c = b\ where \\ & K(i,j) = \int\limits_{\Omega} \varphi_i^* L \varphi_j \ dV & b(i) = \int\limits_{\Omega} \varphi_i \ f \ dV \\ & Key\ insight:\ Calculus\ problem\ of\ solving\ pde\ is\ converted\ to\ linear\ algebra\ problem\ of\ solving\ K\ c = b\ where\ K\ is\ sparse \\ & \end{array}
```



Barnes Hut N-body Simulation

Introduction

- Physical system simulation (time evolution)
 - System consists of bodies
 - "n" is the number of bodies
 - Bodies interact via pair-wise forces
- Many systems can be modeled in these terms
 - Galaxy clusters (gravitational force)
 - Particles (electric force, magnetic force)

Barnes Hut N-body Simulation

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Barnes Hut Idea

- Precise force calculation
 - Requires $O(n^2)$ operations $(O(n^2)$ body pairs)
- Barnes and Hut (1986)
 - Algorithm to approximately compute forces
 - Bodies' initial position & velocity are also approximate
 - Requires only $O(n \log n)$ operations
 - Idea is to "combine" far away bodies
 - Error should be small because force $\sim 1/r^2$

Barnes Hut N-body Simulation

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Barnes Hut Algorithm

- · Set bodies' initial position and velocity
- Iterate over time steps
 - 1. Subdivide space until at most one body per cell
 - · Record this spatial hierarchy in an octree
 - 2. Compute mass and center of mass of each cell
 - 3. Compute force on bodies by traversing octree
 - Stop traversal path when encountering a leaf (body) or an internal node (cell) that is far enough away
 - 4. Update each body's position and velocity

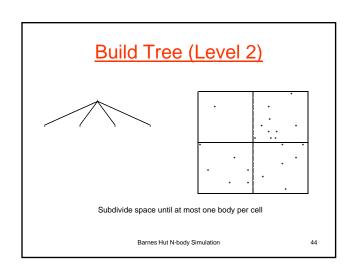
Barnes Hut N-body Simulation

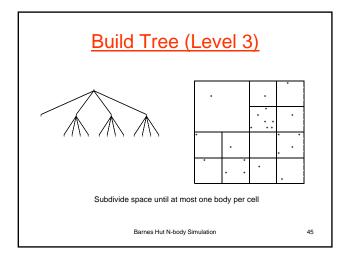
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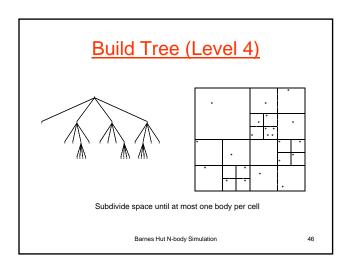
Build Tree (Level 1)

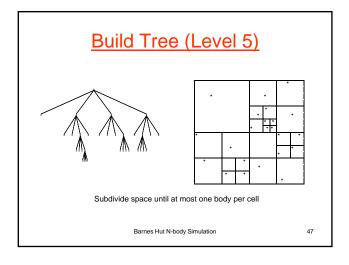
Subdivide space until at most one body per cell

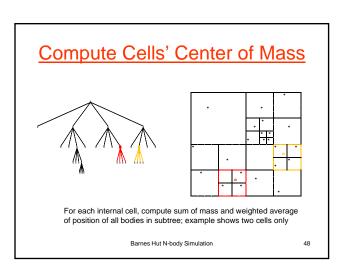
Barnes Hut N-body Simulation

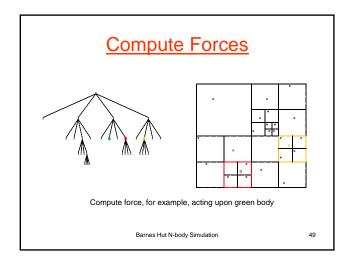


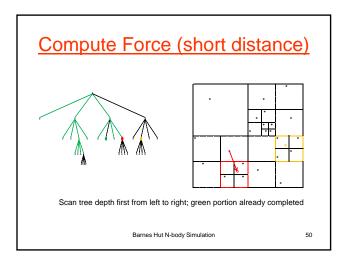


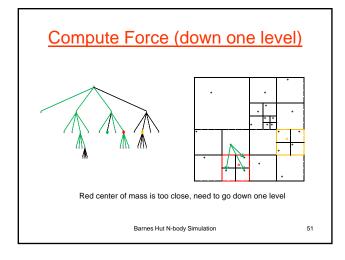


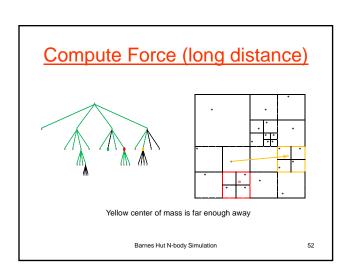


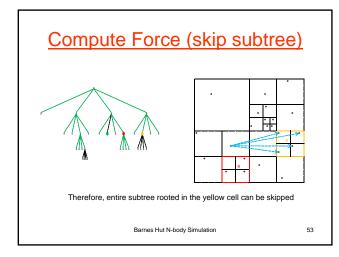












Pseudocode Set bodySet = ... foreach timestep do { Octree octree = new Octree(); foreach Body b in bodySet { octree.Insert(b); } OrderedList cellList = octree.CellsByLevel(); foreach Cell c in cellList { c.Summarize(); } foreach Body b in bodySet { b.ComputeForce(octree); } foreach Body b in bodySet { b.Advance(); } } Barnes Hut N-body Simulation

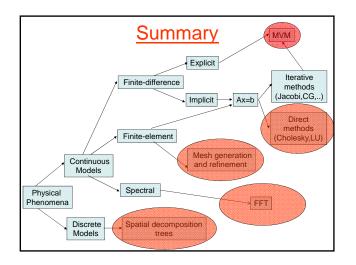
Complexity Set bodySet = ... foreach timestep do { // O(n log n) Octree octree = new Octree(); foreach Body b in bodySet { // O(n log n) octree.Insert(b); OrderedList cellList = octree.CellsByLevel(); foreach Cell c in cellList $\{ // O(n) \}$ c.Summarize(); foreach Body b in bodySet { // O(n log n) b.ComputeForce(octree); foreach Body b in bodySet $\{ // O(n) \}$ b.Advance();

Barnes Hut N-body Simulation

}

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```
Parallelism
Set bodySet = ...
foreach timestep do {
                                // sequential
 Octree octree = new Octree();
 foreach Body b in bodySet { // tree building
    octree.Insert(b);
  OrderedList cellList = octree.CellsByLevel();
  foreach Cell c in cellList { // tree traversal
    c.Summarize();
  foreach Body b in bodySet { // fully parallel
   b.ComputeForce(octree);
  foreach Body b in bodySet { // fully parallel
   b.Advance();
                 Barnes Hut N-body Simulation
                                                  56
```



Summary (contd.)

- Some key computational science algorithms and data structures
 - MVM:
 - explicit finite-difference methods for ode's, iterative linear solvers, finite-element methods
 both dense and sparse matrices
 - stencil computations:
 - finite-difference methods for pde's
 - dense matrices
 - A=LU:

 direct methods for solving linear systems: factorization
 - · usually only dense matrices
 - high-performance factorization codes use MMM as a kernel
 - mesh generation and refinement
 finite-element methods

 - graphs

Summary (contd.)

- Terminology
 - regular algorithms:
 - dense matrix computations like MVM, A=LU, stencil computations
 - parallelism in algorithms is independent of runtime values, so all parallelization decisions can be made at compile-time
 - irregular algorithms:

 - graph computations like mesh generation and refinement
 parallelism in algorithms is dependent on runtime values
 - most parallelization decisions have to be made at runtime during the execution of the algorithm
 - semi-regular algorithms:

 - sparse matrix computations like MVM, A=LU
 parallelization decisions can be made at runtime once matrix is available, but before computation is actually performed
 inspector-executor approach